

# Matrix continued fraction solution of spin-0 and spin-1/2 Feshbach-Villars equation

Bitá Motamedi & Antonio Garcia Vallejo & Zoltán Papp

Department of Physics & Astronomy  
California State University Long Beach

Prague 06/01 - 06/03 2020



# Outline

- 1 Relativistic Quantum Mechanics
- 2 Relativistic QM in Hamiltonian form
- 3 Perspectives and Summary



# Outline

- 1 Relativistic Quantum Mechanics
- 2 Relativistic QM in Hamiltonian form
- 3 Perspectives and Summary



# Quantum Mechanics

- Schrödinger

$$i\hbar\frac{\partial}{\partial t}\psi = -\frac{\hbar^2}{2m}\nabla^2\psi$$



# Quantum Mechanics

- Schrödinger

$$i\hbar\frac{\partial}{\partial t}\psi = -\frac{\hbar^2}{2m}\nabla^2\psi$$

- Klein-Gordon (spin 0 particles)

$$\nabla^2\psi - \frac{1}{c^2}\frac{\partial^2}{\partial t^2}\psi = \frac{mc^2}{\hbar^2}\psi$$



# Quantum Mechanics

- Schrödinger

$$i\hbar \frac{\partial}{\partial t} \psi = -\frac{\hbar^2}{2m} \nabla^2 \psi$$

- Klein-Gordon (spin 0 particles)

$$\nabla^2 \psi - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \psi = \frac{mc^2}{\hbar^2} \psi$$

- Dirac (spin 1/2 particles)

$$i\hbar \gamma^\mu \partial_\mu \psi - mc\psi = 0$$



# Postulates of QM

- 2nd postulate of QM: At any given time the quantum state is described by  $\psi(t)$ .
- 5th postulate of QM: The time evolution is described by

$$i\hbar \frac{d}{dt} |\psi(t)\rangle = H(t) |\psi(t)\rangle$$



# Postulates of QM

- 2nd postulate of QM: At any given time the quantum state is described by  $\psi(t)$ .
- 5th postulate of QM: The time evolution is described by

$$i\hbar \frac{d}{dt} |\psi(t)\rangle = H(t) |\psi(t)\rangle$$

- Klein-Gordon

$$-\hbar^2 \frac{\partial^2}{\partial t^2} \psi = (c^2 p^2 + m^2 c^4) \psi$$





# Postulates of QM

- 2nd postulate of QM: At any given time the quantum state is described by  $\psi(t)$ .
- 5th postulate of QM: The time evolution is described by

$$i\hbar \frac{d}{dt} |\psi(t)\rangle = H(t) |\psi(t)\rangle$$

- Klein-Gordon

$$-\hbar^2 \frac{\partial^2}{\partial t^2} \psi = (c^2 p^2 + m^2 c^4) \psi$$

- The Klein-Gordon equation, one of the basics eq. of relativistic QM, is not a quantum mechanical equation.



# Outline

- 1 Relativistic Quantum Mechanics
- 2 Relativistic QM in Hamiltonian form
- 3 Perspectives and Summary



## Klein-Gordon equation in Hamiltonian form

$$(i\hbar\partial/\partial t - V)^2\psi = (c^2p^2 + (m + S/c^2)^2c^4)\psi$$

$$\psi = \phi + \chi \quad (i\hbar\frac{\partial}{\partial t} - V)\psi = (m + S/c^2)c^2(\phi - \chi) \quad \Psi = \begin{pmatrix} \phi \\ \chi \end{pmatrix}$$

$$i\hbar\frac{\partial}{\partial t}\Psi = H_{FV0}\Psi$$

$$H_{FV}^{(0)} = \begin{pmatrix} 1 & 1 \\ -1 & -1 \end{pmatrix} \left( \frac{p^2}{2m} + U \right) + \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} mc^2 + \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} V$$

Feshbach-Villars (1958)

$$U = S + S^2/2mc^2$$



## Feshbach-Villars equation

$$\phi = \frac{1}{2} \left[ \left( 1 - \frac{1}{mc^2} V \right) \psi + \frac{E}{mc^2} \psi \right]$$

$$\chi = \frac{1}{2} \left[ \left( 1 + \frac{1}{mc^2} V \right) \psi - \frac{E}{mc^2} \psi \right]$$

$$V = 0, E > mc^2, E \simeq mc^2$$

$$\phi = \frac{1}{2} \left[ \psi + \frac{E}{mc^2} \psi \right] \quad \chi = \frac{1}{2} \left[ \psi - \frac{E}{mc^2} \psi \right]$$

$$\phi \rightarrow \psi \quad \chi \rightarrow 0$$



## FV0

$$V = 0, E = mc^2 / \sqrt{1 - v^2/c^2}$$

$$\frac{|\chi|}{|\phi|} = \frac{1 - \sqrt{1 - v^2/c^2}}{1 + \sqrt{1 - v^2/c^2}}$$

as  $v \rightarrow c$   $|\chi|/|\phi| \rightarrow 1$

- if  $E > 0$ ,  $\phi > \chi$   $\chi$  is "shadow" antiparticle component
- if  $E < 0$ ,  $\phi < \chi$   $\phi$  is "shadow" particle component



## Feshbach-Villars for spin-1/2

- $(\gamma_0 P_0 + \gamma_1 P_1 + \gamma_2 P_2 + \gamma_3 P_3 - imc) \psi = 0$
- $(i\hbar\partial/\partial t - V_0)^2 \psi = [\mathbf{p}^2 c^2 + (m + S/c^2)^2 c^4 + c\sigma(\mathbf{p}V_0)]\psi$

$$H_{FV}^{(1/2)} = \begin{pmatrix} 1 & 1 \\ -1 & -1 \end{pmatrix} \left( \frac{p^2}{2m} + U + \frac{\sigma \cdot (pV)}{c} \right) + \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} mc^2 + V \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$



# Hamiltonian formalism

- non-relativistic

$$H = \frac{p^2}{2m} + V$$

- relativistic spin-0

$$H_{FV}^{(0)} = \begin{pmatrix} 1 & 1 \\ -1 & -1 \end{pmatrix} \left( \frac{p^2}{2m} + U \right) + \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} mc^2 + \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} V$$

- relativistic spin-1/2

$$H_{FV}^{(1/2)} = \begin{pmatrix} 1 & 1 \\ -1 & -1 \end{pmatrix} \left( \frac{p^2}{2m} + U + \frac{\sigma \cdot (pV)}{c} \right) + \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} mc^2 + V \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$



## Eigenvalue problem

$$E \begin{pmatrix} |\phi\rangle \\ |\chi\rangle \end{pmatrix} = H_{FV}^{(0)} \begin{pmatrix} |\phi\rangle \\ |\chi\rangle \end{pmatrix}$$

$$H_{FV}^{(0)} = \begin{pmatrix} 1 & 1 \\ -1 & -1 \end{pmatrix} \frac{p^2}{2m} + \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} mc^2 + \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} V$$

- $H_{FV}^{(0)}$  is not hermitian
- yet it has real eigenvalues
- very bad equation: asymptotic coupling by the kinetic energy
- we have to put the asymptotic coupling into the Green's op.





## Lippmann-Schwinger equation

- $H = H_0 + H'$
- $|\psi\rangle = G_0(E)H'|\psi\rangle$        $G(E) = (E - H_0)^{-1}$
- $H' \approx \sum^N |i\rangle\langle i|H'|j\rangle\langle j|$
- $G_0 = |i\rangle\langle i|G_0(E)|j\rangle\langle j|$
- $\underline{\psi} = \underline{G}_0(E)\underline{H}'\underline{\psi}$
- if  $J = E - H_0$  is **tridiagonal** then  $G_0$  is given in terms of **continued fraction**

$$\underline{G} = (\underline{J} - \delta_{iN}\delta_{jN}J_{N,N+1}C_{N+1}J_{N+1,N})^{-1} \quad (1)$$

$$C_{N+1} = (J_{N+1,N+1} - J_{N+1,N+2}C_{N+2}J_{N+2,N+1})^{-1} \quad (2)$$



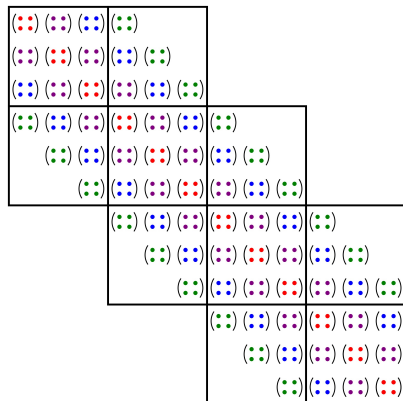
## Coulomb-Sturmian basis

- $S_n(r) = \langle r|n\rangle = \left(\frac{\Gamma(n+1)}{\Gamma(n+2l+2)}\right)^{1/2} e^{-br} (2br)^{l+1} L_n^{2l+1}(2br)$
- 
- $\langle n|E - p^2/(2m) - Z/r|n'\rangle$       tridiagonal
- $\langle n|(E - p^2/(2m) - Z/r)^{-1}|n'\rangle$       continued fraction
- 
- $\langle n|r|n'\rangle$       5-diagonal
- $\langle n|r^2|n'\rangle$       7-diagonal



## FV0 Green's operator for Coulomb plus confining potentials

$$J = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} E - \begin{pmatrix} 1 & 1 \\ -1 & -1 \end{pmatrix} \left( \frac{p^2}{2m} + a_1 r \right) - \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} mc^2 - \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \frac{Z}{r}$$



block - tridiagonal



## FV0 Green's operator for Coulomb plus confining potentials

$$G(E) = \left( \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} E - \begin{pmatrix} 1 & 1 \\ -1 & -1 \end{pmatrix} \left( \frac{p^2}{2m} + a_1 r \right) - \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} mc^2 - \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \frac{Z}{r} \right)^{-1}$$

$$\underline{G} = (\underline{J} - \delta_{iN} \delta_{jN} \underline{J}_{N,N+1} \underline{C}_{N+1} \underline{J}_{N+1,N})^{-1}$$

$$\underline{C}_{N+1} = (\underline{J}_{N+1,N+1} - \underline{J}_{N+1,N+2} \underline{C}_{N+2} \underline{J}_{N+2,N+1})^{-1}$$

continued fraction with block matrices  $\longrightarrow$  matrix continued fraction



## FV1/2 case

$$G(E) = \left( \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} E - \begin{pmatrix} 1 & 1 \\ -1 & -1 \end{pmatrix} \left( \frac{p^2}{2m} + \frac{\sigma \hbar Z}{c} \frac{1}{r^2} \right) - \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} mc^2 - \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \frac{Z}{r} \right)^{-1}$$

$$l \rightarrow -1/2 + \sqrt{1/4 + (j + 1/2)^2 \pm \sqrt{(j + 1/2)^2 - Z^2 \alpha^2}}$$

perfect agreements with Klein-Gordon and Dirac results !!!



# Outline

- 1 Relativistic Quantum Mechanics
- 2 Relativistic QM in Hamiltonian form
- 3 Perspectives and Summary



# Feshbach-Villars formalism

- relativistic quantum mechanical equation
- multi-component wave function
- linear time & quadratic spatial derivative
- non-hermitian Hamiltonian
  - $\bar{\Omega} = \tau_3 \Omega^\dagger \tau_3$       FV adjoint
  - $\bar{\Omega} = \Omega$               FV hermitian  $\rightarrow$  real eigenvalues
- explicit particle - antiparticle components
- interesting formalism, worth studying

